

Quantum Foundations Talking Points 31/5/19

§§3.2–3.5 of Notes

Pairs of Qubits

- General joint state:

$$\alpha_{00} |0\rangle \otimes |0\rangle + \alpha_{01} |0\rangle \otimes |1\rangle + \alpha_{10} |1\rangle \otimes |0\rangle + \alpha_{11} |1\rangle \otimes |1\rangle, \quad \sum_{j,k} |\alpha_{jk}|^2 = 1$$

$$= \alpha_{00} |00\rangle + \alpha_{01} |01\rangle + \alpha_{10} |10\rangle + \alpha_{11} |11\rangle.$$
- physical realisations (polarised light, energy levels, spin)
- may be physically separated

Bell States

- $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$, $|\Phi^-\rangle = \frac{1}{\sqrt{2}}(|00\rangle - |11\rangle)$, $|\Psi^+\rangle = \frac{1}{\sqrt{2}}(|01\rangle + |10\rangle)$, $|\Psi^-\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$.
- Not product states, are “maximally” entangled
- orthonormal basis

Invariance Properties

- Invariance property for $|v\rangle = U|0\rangle$, $|v^\perp\rangle = U|1\rangle$, real unitary $U = \begin{bmatrix} \alpha & -\beta \\ \beta & \alpha \end{bmatrix}$.
Then $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) = \frac{1}{\sqrt{2}}(|vv\rangle + |v^\perp v^\perp\rangle)$
- $|\Psi^-\rangle$ has full invariance
- Measure Alice’s qubit w.r.t. any $|v\rangle$, Bob’s qubit gives same result with $|v\rangle$.
Synchronised coins.
- Can be classically explained, but not so other correlations.

CNOT

- unitary operating on pair of qubits, both need to be present
-

$$\begin{array}{cccc} |00\rangle \mapsto |00\rangle & |01\rangle \mapsto |01\rangle & |10\rangle \mapsto |11\rangle & |11\rangle \mapsto |10\rangle \\ \begin{array}{c} |0\rangle \xrightarrow{\bullet} |0\rangle \\ |0\rangle \xrightarrow{\oplus} |0\rangle \end{array} & \begin{array}{c} |0\rangle \xrightarrow{\bullet} |0\rangle \\ |1\rangle \xrightarrow{\oplus} |1\rangle \end{array} & \begin{array}{c} |1\rangle \xrightarrow{\bullet} |1\rangle \\ |0\rangle \xrightarrow{\oplus} |1\rangle \end{array} & \begin{array}{c} |1\rangle \xrightarrow{\bullet} |1\rangle \\ |1\rangle \xrightarrow{\oplus} |0\rangle \end{array} \end{array}$$

- control and target
- produces entanglement

$$\left. \begin{array}{c} \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\ |0\rangle \xrightarrow{\oplus} \end{array} \right\} |\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

- universal property

- matrix $\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$.

Preparing Bell States

$$\left. \begin{array}{c} |0\rangle \xrightarrow{[H]} |0\rangle \\ |0\rangle \xrightarrow{\oplus} \end{array} \right\} |\Phi^+\rangle$$

$$\left. \begin{array}{c} |0\rangle \xrightarrow{[H]} |1\rangle \\ |1\rangle \xrightarrow{\oplus} \end{array} \right\} |\Psi^+\rangle$$

$$\left. \begin{array}{c} |1\rangle \xrightarrow{[H]} |1\rangle \\ |0\rangle \xrightarrow{\oplus} \end{array} \right\} |\Phi^-\rangle$$

$$\left. \begin{array}{c} |1\rangle \xrightarrow{[H]} |0\rangle \\ |1\rangle \xrightarrow{\oplus} \end{array} \right\} |\Psi^-\rangle$$

Tensor Products of Operators

- Example: $H \otimes I$
- In general: $S \otimes T$, $|\psi_1\rangle \xrightarrow{[S]} S|\psi_1\rangle$
 $|\psi_2\rangle \xrightarrow{[T]} T|\psi_2\rangle$

$$(S \otimes T)(|\psi_1\rangle \otimes |\psi_2\rangle) = S|\psi_1\rangle \otimes T|\psi_2\rangle$$

$$\bullet S \otimes T = \begin{bmatrix} S_{11}T_{11} & S_{11}T_{12} & S_{12}T_{11} & S_{12}T_{12} \\ S_{11}T_{21} & S_{11}T_{22} & S_{12}T_{21} & S_{12}T_{22} \\ S_{21}T_{11} & S_{21}T_{12} & S_{22}T_{11} & S_{22}T_{12} \\ S_{21}T_{21} & S_{21}T_{22} & S_{22}T_{21} & S_{22}T_{22} \end{bmatrix} = \begin{bmatrix} S_{11}T & S_{12}T \\ S_{21}T & S_{22}T \end{bmatrix}.$$

- CNOT is not a tensor product (tensor products don't entangle, or from matrix representation)
- Tensoring preserves unitary and hermitian properties

Measuring Entangled States — Partial Measurement Rule

- Single qubit measured w.r.t. $\{|v\rangle, |v^\perp\rangle\}$:
 $\alpha|v\rangle + \beta|v^\perp\rangle \rightarrow \alpha|v\rangle + \beta/\sqrt{\alpha^2 + \beta^2}|v^\perp\rangle \rightarrow \frac{\alpha}{\sqrt{\alpha^2 + \beta^2}}|v\rangle \sim |v\rangle$
Similarly for outcome $|v^\perp\rangle$.
- Measure first qubit
 $a|vw\rangle + b|vw^\perp\rangle + c|v^\perp w\rangle + d|v^\perp w^\perp\rangle = |v\rangle \otimes (a|w\rangle + b|w^\perp\rangle) + |v^\perp\rangle \otimes (c|w\rangle + d|w^\perp\rangle)$
 $\rightarrow |v\rangle \otimes (a|w\rangle + b|w^\perp\rangle) + \frac{a}{\sqrt{a^2 + b^2}}|w\rangle + \frac{b}{\sqrt{a^2 + b^2}}|w^\perp\rangle$ with probability $|a|^2 + |b|^2$,

Similarly for outcome $|v^\perp\rangle$.

- Measure second qubit
 $|v\rangle \otimes |w\rangle$ with probability $(|a|^2 + |b|^2) \frac{|a|^2}{|a|^2 + |b|^2} = |a|^2$
 $|v\rangle \otimes |w^\perp\rangle$ with probability $(|a|^2 + |b|^2) \frac{|b|^2}{|a|^2 + |b|^2} = |b|^2$
Similarly $|v^\perp\rangle \otimes |w\rangle, |v^\perp\rangle \otimes |w^\perp\rangle$.

- In computational basis, measurement operator is $\begin{bmatrix} \lambda_{00} & 0 & 0 & 0 \\ 0 & \lambda_{01} & 0 & 0 \\ 0 & 0 & \lambda_{10} & 0 \\ 0 & 0 & 0 & \lambda_{11} \end{bmatrix}$ provided λ_{ij} all distinct.
- Bell $|\Phi^+\rangle$ in computational basis, outcome $|00\rangle$ or $|11\rangle$ each with probability $1/2$.

Measuring w.r.t. Bell Basis

- $|00\rangle, |01\rangle, |10\rangle, |11\rangle \xrightarrow{\left\{ \begin{array}{c} H \\ \text{---} \\ \oplus \end{array} \right\}} |\Phi^+\rangle, |\Psi^+\rangle, |\Phi^-\rangle, |\Psi^-\rangle$, hence
 $|\Phi^+\rangle, |\Psi^+\rangle, |\Phi^-\rangle, |\Psi^-\rangle \xrightarrow{\left\{ \begin{array}{c} \oplus \\ \text{---} \\ H \end{array} \right\}} |00\rangle, |01\rangle, |10\rangle, |11\rangle$. Explain.
- Hence physically realise measurement in the Bell basis

Superdense Coding Comments

- “normally” 1 qubit codes up 1 bit
- A & B share Bell pair, A transforms her qubit, sends to Bob, he extracts **2** bits of info.

Superdense Coding Method

- A & B share a pair of qubits in Bell state $|\Phi^+\rangle$
- To send 00, 01, 10 or 11 to B, A applies $I \otimes I$, $\sigma_3 \otimes I$, $\sigma_1 \otimes I$, or $i\sigma_2 \otimes I$ respectively to $|\Phi^+\rangle$.

$$|\Phi^+\rangle \left\{ \begin{array}{c} [I] \\ \text{---} \end{array} \right\} |\Phi^+\rangle \quad |\Phi^+\rangle \left\{ \begin{array}{c} [\sigma_3] \\ \text{---} \end{array} \right\} |\Phi^-\rangle$$

$$|\Phi^+\rangle \left\{ \begin{array}{c} \text{---} \boxed{\sigma_1} \text{---} \\ \text{---} \end{array} \right\} |\Psi^+\rangle \quad |\Phi^+\rangle \left\{ \begin{array}{c} \text{---} \boxed{i\sigma_2} \text{---} \\ \text{---} \end{array} \right\} |\Psi^-\rangle$$

- A sends her transformed qubit to B
- B measures the pair in the Bell basis

Teleportation Comments

- Alice has qubit $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$, wants to send state to Bob
- No electronic connection, just mail postal service
- Share a $|\Phi^+\rangle$ state
- A can send 2 classical bits to Bob, he can then reconstruct $|\psi\rangle$
- *Impressive!!*

Teleportation Method

- $|\psi_0\rangle = |\psi\rangle \otimes |\Phi^+\rangle = \frac{1}{\sqrt{2}}(\alpha|000\rangle + \alpha|011\rangle + \beta|100\rangle + \beta|111\rangle)$
 $= \frac{1}{2}\left(|\Phi^+\rangle(\alpha|0\rangle + \beta|1\rangle) + |\Phi^-\rangle(\alpha|0\rangle - \beta|1\rangle) + |\Psi^+\rangle(\beta|0\rangle + \alpha|1\rangle) + |\Psi^-\rangle(-\beta|0\rangle + \alpha|1\rangle)\right).$
- Alice measures in Bell basis, gets one of 4 results, mails which result to Bob
- Bob applies gates $I, \sigma_3, \sigma_1, i\sigma_2 = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$ resp. Always gets $\alpha|0\rangle + \beta|1\rangle$. *Impressive!!*
- Note: Bob not entangled until Alice measures in Bell basis.